









Master I.M.A.G.2E

Imagerie et Modélisation Astrophysique et Géophysique, Espace et Environnement

Master 2– Stellar Physics

Dr. Philippe Stee Observatoire de la Côte d'Azur – CNRS Philippe.Stee@oca.eu

Slides mainly from P. Crowther

Recommended Texts

- Bohm-Vitense, Introduction to Stellar Astrophysics: Vol 2, ISBN 0521348706 Essential (full derivation of important equations)
- Gray, Observations and analysis of stellar photospheres, ISBN 0521408687, <u>Recommended</u> (good background reading, but complex maths)

cgs (cm/gram/second) units

Fundamental:

Gravitational constant, G=6.679x10⁻⁸ cm³/g/s² Stefan-Boltzmann: σ =5.6705x10⁻⁵ erg/cm²/s/K⁴ Speed of light, c=2.99792x10¹⁰cm, Electron mass: m_e=9.109x10⁻²⁸ g Planck constant h=6.626x10⁻²⁷ erg s Electron charge: 4.803x10⁻¹⁰ e.s.u. Boltzmann const: k=1.380x10⁻¹⁶ erg/s or 8.617x10⁻⁵ eV/K Gas constant R=8.314x10⁷ erg/mol K

Solar

Radius R_{\odot} =6.955x10¹⁰ cm, Luminosity L_{\odot}= 3.845x10³³ erg/s, Mass M_{\odot}=1.989x10³³ g

Astronomical:

Parsec, pc: 3.085x10¹⁸cm

Astronomical Unit, AU: 1.496x10¹³ cm

Miscellaneous

Energy of 1ev=1.602x10⁻²⁶ erg

Color/Temperature Relation



Line Spectrum (e.g. Sun)



Fraunhofer (1814)

Continuous Spectrum (e.g. Vega)



What is a stellar atmosphere?

- Thin, tenuous transition zone between (invisible) stellar interior and (essentially vacuum) exterior.
- The `photosphere' is the visible disk, whilst the `atmosphere' also includes coronae and winds.
- In contrast with the interior, where convection dominates, the energy transport mechanism of the atmosphere is radiation.
- Stellar atmospheres are primarily characterized by two parameters: (Teff, log g)

$(T_{eff}, \log g)$

- Effective temperature (K), is defined by L= $4\pi R^2 \sigma T_{eff}^4$ related to *ionization*.
- Surface gravity (cm/s²), $g = GM/R^2$, related to *pressure*.
- The Sun has T_{eff}=5777K, log g=4.44 its atmosphere is only a few hundred km deep, <0.1% of the stellar radius.
- Solar atmosphere most easily studied during total eclipse (lunar limb occults at rate of 0.5 arcsec/s or 300km/s at 1AU, so brightness variation during last second before totality provides physical size).
- A red giant has log g~1 (extended atmosphere) whilst a white dwarf has log g~8 (no atmosphere).

Spectral Types

M-K (Morgan-Keenan) classification scheme orders stars via "OBAFGKM" spectral types using ratios of line strength. O-types have the bluest B-V & highest Teff's. Each are subdivided into (up to) ten divisions – e.g. O2 .. O9, B0, B1 .. B9, A0, A1 .. etc

Table 15.1. MK spectral classes.			
MK spectral class	Class characteristics		
0	Hot stars with He II absorption		
В	He I absorption; H developing later		
Α	Very strong H, decreasing later; Ca II increasing		
F	Ca II stronger; H weaker; metals developing		
G	Ca II strong; Fe and other metals strong; H weaker		
K	Strong metallic lines; CH and CN bands developing		
М	Very red; TiO bands developing strongly		



Figure 8.9 The dependence of spectral line strengths on temperature.

Stellar Spectra



Luminosity Class

Luminosity class information is often added from line widths: V (dwarfs), III (giants), I (supergiants), e.g. B0III is an early B giant, which is itself an `early-type' star (along with O and A-types), whilst cooler stars are `late-type'

Class	Line pairs for class	Class	Line pairs for luminosity
05 ⇔ 09	4471 He 1/4541 He II	09 ⇔ B3	4116-21 (Si IV, He I)/4144 He I
B0 ⇔ B1	4552 Si III/4089 Si IV	B0 ⇔ B3	3995 N 11/4009 He 11
B2 ⇔ B8	4128-30 Si 11/4121 He I	B1 ⇔ A5	Balmer line wings
B8 ⇔ A2	4471 He 1/4481 Mg II 4026 He 1/3934 Ca II	A3 ⇔ F0	4416/4481 Mg II
A2 ⇔ F5	4030-34 Mn 1/4128-32 4300 CH/4385	F0 ⇔ F8	4172/4226 Ca I
F2 ⇔ K	4300 (G band)/4340 Hy	F2 ⇔ K5	4045-63 Fe 1/4077 Sr II
F5 ⇔ G5	4045 Fe 1/4101 Hδ		4226 Ca 1/4077 Sr II
	4226 Ca 1/4340 Hy	G5 ⇔ M	Discontinuity near 4215
G5 ⇔ K0	4144 Fe 1/4101 H8	K3 ⇔ M	4215/4260, Ca I increasing
K0 ⇔ K5	4226 Ca 1/4325 4290/4300		a double shifts of a subich o

Luminosity classes

- la bright supergiant
- Ib Supergiant
- II bright giant
- III giant
- IV subgiant
- V main-sequence star

Example Luminosity Classes

Our Sun: G2 star on the Main Sequence: G2V

Polaris: G2 star with Supergiant luminosity: G2lb

Stellar Spectra

Class	Colour	Temp. (x10 ³ K)	Spectral lines	Examples
0	Blue-violet	28 – 50	lonised atoms	z Pup ¹ , d Ori ²
В	Blue-white	10 – 28	He, some H	a Vir ³ , b Ori ⁴
A	White	7.5 – 10	Strong H, some ionised metals	a Cma ⁵ , a Lyr ⁶
F	Yellow-white	6 – 7.5	H and Ca ⁿ⁺ , Fe ⁿ⁺ .	a Car ⁷ , a Cmi ⁸
G	Yellow	5 – 6	Ca ⁿ⁺ , other ionised and neutral metals	Sun, a Aur ⁹
K	Orange	3.5 – 5	Neutral metals	a Boo ¹⁰ , a Tau ¹¹
М	Red-orange	2.5 – 3.5	TiO and Ca	a Sco ¹² , a Ori ¹³

Common names: ¹Naos, ²Mintaka, ³Spica, ⁴Rigel, ⁵Sirius, ⁶Vega, ⁷Canopus, ⁸Procyon, ⁹Capella, ¹⁰Arcturus, ¹¹Aldebaran, ¹²Antares

H-R Diagram for a number of the brightest and nearest stars



The H-R Diagram

- Points to note:
 - The narrow band of stars scattered close to the solid line.
 - Most stars occur along this band an indication that this is where stars spend most of their lives. For this reason, it is known as the *Main Sequence*.

The H-R Diagram

- Other regions to note are stars of high luminosity but low temperature (indicating they are large – hence the term *red giant*) and stars of high temperature but low luminosity (indicating small diameters, hence *white dwarf*)
- As we shall see, the H-R diagram is extremely useful in many aspects of stellar physics

Hertzsprung-Russell Diagram





Magnitude scale

- We measure the flux F from astronomical objects via a logarithmic magnitude scale (like the eye).
- A star which is 5 mags more negative than another implies it is a factor of 100 times brighter
- Often involves broad-band filters (UBV at optical wavelengths), e.g. for V-band

 $m_v - m_0 = -2.5 \log(F_v/F_0)$

 Vega (A0V) defines the photometric `zero point' m₀ at all wavelengths (U=B=V=0.0 mag etc).

Table 15.6. Flux calibration for an A0 V star.		
Symbol	Flux (erg cm ⁻² s ⁻¹ Å ⁻¹)	$\lambda_0 (\mu m)$
U	4.22×10^{-9}	0.36
B	6.40×10^{-9}	0.44
V	3.75×10^{-9}	0.55
R	1.75×10^{-9}	0.71
I	8.4×10^{-10}	0.97

IR magnitudes

Vega (or Sirius) also defines `zero point' in Infrared. IR filters chosen to match transmission through Earth's atmosphere which is low at some wavelengths even from high, dry sites such as Mauna Kea:

Filter name	λ_{iso}^{b} (μ m)	$\Delta \lambda^c$ (μ m)	$(W m^{-2} \mu m^{-1})$
V	0.5556 ^d		3.44×10^{-8}
J	1.215	0.26	3.31×10^{-9}
Н	1.654	0.29	1.15×10^{-9}
Ks	2.157	0.32	4.30×10^{-10}
K	2.179	0.41	4.14×10^{-10}
L	3.547	0.57	6.59×10^{-11}
L'	3.761	0.65	5.26×10^{-11}
М	4.769	0.45	2.11×10^{-11}
8.7	8.756	1.2	1.96×10^{-12}
N	10.472	5.19	9.63×10^{-13}
11.7	11.653	1.2	6.31×10^{-13}
Q	20.130	7.8	7.18×10^{-14}



Colour index

We can define a colour index as the difference between filters relative to Vega e.g. $B-V = m_B-m_{V,}$ such that stars bluer than A0 have a –ve B-V colour and red stars a +ve colour. The Sun has B-V=+0.65 mag.



More on magnitudes

• We define the absolute (visual) magnitude (M_V) as the apparent (visual) magnitude of a star of m_V lying at a distance of 10pc:

 $M_V = m_V - 5\log(d/pc) + 5 - A_V$.

- For the Sun (d=4.85 10⁻⁶ pc), m_V =-26.75 and M_V =+4.82 mag. The distance modulus is M_V -m_v.
- Interstellar extinction requires the A_V term, where $A_V{\sim}3.1$ E(B-V) for most sight lines with

 $E(B-V)=B-V - (B-V)_{o}$

i.e. the difference between the observed and intrinsic colour.

Interstellar Reddening

One also needs to correct color indices for interstellar reddening. As the light propagates through interstellar dust, the blue light is scattered preferentially making objects appear to be redder than they actually are...



Interstellar Extinction

Extinction is MUCH higher at shorter wavelengths, so IR observations of e.g. Milky Way disk probe much further. The extinction to the Galactic Centre (8kpc) is approx A_V =30 mag (at 5500A) versus A_K =3 mag (at 2micron).



Bolometric Flux

The bolometric flux from a star (erg cm⁻² s⁻¹) received at the top of the Earth's atmosphere is the integral of the spectral flux (measured at a frequency Fv) over all wavelengths:

$$F_{Bol} = \int_0^\infty F_{\nu} d\nu$$

The luminosity (erg/s) is the bolometric flux from the star integrated over a full sphere (at distance d):

$$L = 4\pi d^2 F_{Bol}$$

Since the Earth's atmosphere is opaque to UV and some IR radiation one cannot always directly measure the bolometric flux.

Bolometric Corrections

One can calculate bolometric corrections (BC) from atmospheric models to correct measured fluxes (usually in the V band) for the total flux. Usually expressed in magnitudes:

BC=
$$M_{bol} - M_V$$
 with $M_{bol} = 4.74 - 2.5 \log(L/L_{\odot})$

BC=-0.08 mag for the Sun is a small correction since it emits most radiation in the visual. Hot OB stars have very negative BC's, since most of the energy is emitted in the UV, as are cool M stars with most energy emitted in the IR.

Bolometric Corrections

Bolometric corrections can be estimated from intrinsic colours (B-V)o as shown here for dwarfs:





Or from the Spectral Type, using a Teff-Spectral Type calibration (e.g. Astrophysical Quantities see next slide)

Properties of Main-Sequence Stars

Table 15.7. Calibration of MK spectral types.							
Sp	M(V)	B - V	U - B	V - R	R-I	T _{eff}	BC
MAI	N SEQUEN	ICE, V					
05	-5.7	-0.33	-1.19	-0.15	-0.32	42 000	-4.40
09	-4.5	-0.31	-1.12	-0.15	-0.32	34 000	-3.33
B0	-4.0	-0.30	-1.08	-0.13	-0.29	30 000	-3.16
B2	-2.45	-0.24	-0.84	-0.10	-0.22	20 900	-2.35
B5	-1.2	-0.17	-0.58	-0.06	-0.16	15 200	-1.46
B 8	-0.25	-0.11	-0.34	-0.02	-0.10	11 400	-0.80
A0	+0.65	-0.02	-0.02	0.02	-0.02	.9 790	-0.30
A2	+1.3	+0.05	+0.05	0.08	0.01	9 000	-0.20
A5	+1.95	+0.15	+0.10	0.16	0.06	8 180	-0.15
F0	+2.7	+0.30	+0.03	0.30	0.17	7 300	-0.09
F2	+3.6	+0.35	0.00	0.35	0.20	7 000	-0.11
F5	+3.5	+0.44	-0.02	0.40	0.24	6 6 5 0	-0.14
F8	+4.0	+0.52	+0.02	0.47	0.29	6250	-0.16
G0	+4.4	+0.58	+0.06	0.50	0.31	5940	-0.18
G2	+4.7	+0.63	+0.12	0.53	0.33	5790	-0.20
G5	+5.1	+0.68	+0.20	0.54	0.35	5 560	-0.21
G8	+5.5	+0.74	+0.30	0.58	0.38	5310	-0.40
K0	+5.9	+0.81	+0.45	0.64	0.42	5 1 5 0	-0.31
K2	+6.4	+0.91	+0.64	0.74	0.48	4830	-0.42
K5	+7.35	+1.15	+1.08	0.99	0.63	4410	-0.72
M0	+8.8	+1.40	+1.22	1.28	0.91	3 8 4 0	-1.38
M2	+9.9	+1.49	+1.18	1.50	1.19	3 5 2 0	-1.89
M5	+12.3	+1.64	+1.24	1.80	1.67	3 170	-2.73

Example

- <u>A BOV in the LMC (distance 50kpc) has V=13.0 mag and B-</u> V=-0.20 mag. What is its bolometric luminosity, relative to the Sun?</u>
- 1. From the Table, E(B-V)=(B-V)-(B-V)o=0.1 mag for a B0 dwarf, so $A_V=0.3$ mag (the star is lightly reddened).
- 2. Using $M_v=m_v-5\log(d/pc)+5-A_{v,}$ we obtain $M_v=13-23.5+5-0.3=-5.8$ mag.
- 3. From the Table, BC(B0V)=-3.16, so -3.16= M_{bol} -(-5.8), hence M_{bol} =-9 mag
- 4. From M_{bol} =4.74-2.5xlog(L/L_{\odot}), we obtain log(L/L_{\odot})=5.5, some 300,000 times brighter than the Sun after integrating over all λ

H-R diagram

- 90% of stars are on the main sequence and obey the massluminosity dependence L ~ M^{3.5}
- Stars on the main sequence generate energy due to nuclear fusion of hydrogen
- In the end of their lives stars move to the upper right corner of the H-R diagram



Check this hypothesis

- Mass should be most important parameter
- It determines the pressure in the star center and the central temperature
- It determines the surface temperature

$$\frac{L \propto M^{3.5}}{L_{sun}} = \left(\frac{M}{M_{sun}}\right)^{3.5}$$

Gravity Holds a Star Together

Stars are held together by gravity. Gravity tries to compress everything to the center. What holds an ordinary star up and prevents total collapse is thermal and radiation pressure. The thermal and radiation pressure tries to expand the star layers outward to infinity.

- 1. Newton's gravitation law
- 2. Hydrostatic equilibrium
- 3. Equation of state
- 4. Energy transport

Mass determines all star's properties



Amount of hydrogen fuel Lifetime = Rate of energy loss Lifetime T ~ M/L ~ $1/M^{3.5-1} = 1/M^{2.5}$ M = 4M_X; $\frac{T}{T_{sun}} = \left(\frac{M_{sun}}{M}\right)^{2.5} = \frac{1}{32}$ $T \sim 3 \times 10^8$ years star mass (solar masses) time (years) 3 million

60

Spectral type 03 07

30	11 million	07
10	32 million	B4
3	370 million	A5
1.5	3 billion	F5
1	10 billion	G2 (Sun)
0.1	1000's billions	M7

TABLE 9-2 Main-Sequence Stars

Spectral Type	Mass (Sun = 1)	Luminosity (Sun = 1)	Years on Main Sequence
05	40	405,000	$1 imes 10^6$
B0	15	13,000	$11 imes10^6$
AO	3.5	80	440×10^{6}
FO	1.7	6.4	$3 imes10^9$
GO	1.1	1.4	$8 imes10^9$
KO	0.8	0.46	17×10^9
MO	0.5	0.08	$56 imes 10^9$

How to explain the cutoff at masses > 100 M_{sun} and < 0.08 M_{sun}



Maximum Masses of Main-Sequence Stars

 $M_{max} \sim 50$ - 100 solar masses

a) More massive clouds fragment into
b) Very massive stars lose
smaller pieces during star formation.
mass in strong stellar winds



Example: η Carinae: Binary system of a 60 M_{sun} and 70 M_{sun} star. Dramatic mass loss; major eruption in 1843 created double lobes.

High-mass cutoff at M ~ 100 M_{sun}

Too massive and luminous stars throw off their outer layers due to radiation pressure

Eta Carinae





PISTOI NEBUIA and MASSIVE Star HST • NICMO PRC97-33 • ST ScI OPO • D. Figer (UCLA) and NASA

Minimum Mass of Main-Sequence Stars

 M_{min} = 0.08 M_{sun}



At masses below 0.08 M_{sun}, stellar progenitors do not get hot enough to ignite thermonuclear fusion.

→ Brown Dwarfs

Low-mass cutoff of the main sequence: $M \sim 0.08 M_{sun}$

Brown dwarfs: temperature is too low to ignite nuclear fusion



Gliese 229B: only 0.02 M_{sun}

Conclusion

Based on this evidence, we conclude:

Stars spend most of their lives as main sequence stars.

During its lifetime, the surface temperature and luminosity stays almost constant.

Something else could happen in the star birth process.

Something else could happen in the star death process.

The star's mass determines what the temperature and luminosity is during the star's main sequence lifetime.

More mass -> hotter. More mass -> more luminous.

Also, more mass -> bigger.



Radiation Terms

Black body radiation (Planck function) Effective Temperature (Stefan-Boltzmann) Specific and mean Intensity

The Black Body

Imagine a box which is completely closed except for a small hole. Any light entering the box will have a very small likely hood of escaping & will eventually be absorbed by the gas or walls. For constant temperature walls, this is in thermodynamic equilibrium.



If this box is heated the walls will emit photons, filling the inside with radiation. A small fraction of the radiation will leak out of the hole, but so little that the gas within it remains in equilibrium. The emitted radiation is that of a black-body. Stars share properties of the black-body emitter, in the sense that a negligibly small fraction of the radiation escapes from each.

Stars do differ from black bodies

The observed flux distributions of real stars deviate from black body curves, as indicated here for the UBV colours of dwarfs and supergiants. This difference is due to sources of continuous and line opacity in the stellar photospheres.

